

**Higher Spin Theory
and Holography-4
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**Ordinary-derivative approach to
long conformal fields**

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Plan

- 1) **Introduction - Review-Results**
- 2) **Modified Lorentz and de Donder gauges
and (global) BRST Lagrangian for conformal fields**
- 3) **Computation of partition functions
of long and partial-short conformal fields**

Totally symmetric conformal field in $R^{d-1,1}$

$$\phi^s = \phi^{a_1 \dots a_s}$$

$$\Delta = \frac{d}{2} - \kappa$$

$$\mathcal{L} = \phi^s \square^\kappa \phi^s$$

Three arbitrary integer labels

s, κ , d

Short conformal field

$$\kappa = s + \frac{d-4}{2}, \quad d - \text{even}$$

Partial-short conformal field

$$\kappa = s + \frac{d-4}{2} - t, \quad d - \text{even}$$

$$t = 1, 2, \dots, s - 1$$

Long conformal field

$$\kappa > s + \frac{d-4}{2}, \quad \kappa - \text{integer}$$

In **AdS/CFT**

Conformal field in $R^{d-1,1}$ with

$$\Delta = \frac{d}{2} - \kappa$$

is dual to -non-normalizable modes of field

in AdS_{d+1} with

$$E_0 = \frac{d}{2} + \kappa$$

Short conformal spin-2 field in $R^{3,1}$ (Weyl Gravity)

$$Z = \frac{(D_1)^3}{(D_2)^2}$$

$$n^{\text{DoF}} = 9 \times 2 - 4 \times 3 = 6$$

$$D_n \equiv \sqrt{\det(-\square)}$$

for rank- n traceless tensor field

Fradkin, Tseytlin 1981

also n^{DoF} by Dirac method

Nieuwenhuizen, Lee 1982

Short conformal arbitrary spin- s field in $R^{3,1}$ (Fradkin-Tseytlin fields)

$$\mathcal{L} = \phi^S \square^S \phi^S$$

$$Z = \frac{(D_{s-1})^{s+1}}{(D_s)^s}$$

$$\begin{aligned} n^{\text{DoF}} &= s \times (s+1)^2 - (s+1) \times s^2 \\ &= s(s+1) \end{aligned}$$

Fradkin, Tseytlin 1985

Short conformal arbitrary spin- s field in $R^{d-1,1}$

$$Z = \frac{(D_{s-1})^{\nu_s+1}}{(D_s)^{\nu_s}} \quad \nu_s \equiv s + \frac{d-4}{2}$$

$$\begin{aligned} n^{DoF} &= \nu_s n_s^{\text{so}(d)} - (\nu_s + 1) n_{s-1}^{\text{so}(d)} \\ &= \frac{1}{2}(d-3)(2s+d-2)(2s+d-4) \frac{(s+d-4)!}{(d-2)!s!} \end{aligned}$$

Tseytlin 2013

also n^{DoF} by light-cone gauge method

RRM 2007

Partial-short maximal-dept conformal spin- s field in $R^{3,1}$

$$Z = \frac{(D_0)^{s+1}}{(D_s)}$$

$$n^{DoF} = s(s + 1)$$

$$\kappa = s + \frac{d - 4}{2} - t$$

$$\kappa = 1, \quad d = 4, \quad t = s - 1$$

Beccaria, Tseytlin 2015

Long conformal arbitrary spin- s field in $R^{d-1,1}$

$$Z = \frac{1}{(D_s)^\kappa}$$

$$n^{DoF} = \kappa n_s^{\text{so}(d)}$$

$$n_s^{\text{so}(d)} \equiv (2s + d - 2) \frac{(s + d - 3)!}{(d - 2)!s!}$$

$$n^{\text{DoF}} = \kappa \times (\text{DoF massive field in } (d+1) \text{ dimensions})$$

Partial-short conformal spin- s field in $R^{d-1,1}$

$$Z = \frac{(D_{s-1-t})^{s+\frac{d-2}{2}}}{(D_s)^\kappa}$$

$$\kappa = s + \frac{d-4}{2} - t$$

$$n^{DoF} = \frac{(2s+d-2)(2s+d-4-2t)(s+d-4)!}{2(d-2)!s!} \\ \times \left(s+d-3 - s \frac{(s-t)_t}{(s-t+d-3)_t} \right),$$

Arbitrary values of κ, s, d

$$\mathcal{L} = \phi^S \square^\kappa \phi^S$$

ϕ^S traceless rank- s tensor field of $so(d-1, 1)$

higher-derivative Lagrangian

Vasiliev 2009

spin-1

$$\mathcal{L} = -\frac{1}{4}F^{ab}F^{ab}$$

$$F^{ab} = \partial^a\phi^b - \partial^b\phi^a$$

$$\mathcal{L} = \frac{1}{2}\phi^a\Box\phi^a + \frac{1}{2}L^2$$

$$L = \partial^a\phi^a$$

$$\Box\phi^a = 0 \quad L = 0$$

Apply Faddeev-Popov procedure

Step 1. Introduce

Faddeev-Popov fields

\bar{c} c

Nakanishi-Lautrup field

b

Step 2.

$$\mathcal{L}^{BRST} = \mathcal{L} + \mathcal{L}_{\text{g.fix}}$$

$$\mathcal{L}_{\text{g.fix}} = -\mathbf{bL} + \bar{\mathbf{c}}\square\mathbf{c} + \frac{1}{2}\xi\mathbf{b}^2$$

$$\xi = 0$$

Landau gauge

$$\xi = 1$$

Feynman gauge

BRST

$$s\phi^a = \partial^a c$$

$$sc = 0$$

$$s\bar{c} = b$$

$$sb = 0$$

anti-BRST

$$\bar{s}\phi^a = \partial^a \bar{c}$$

$$\bar{s}c = -b$$

$$\bar{s}\bar{c} = 0$$

$$\bar{s}b = 0$$

$$s^2 = 0$$

$$\bar{s}^2 = 0$$

$$s\bar{s} + \bar{s}s = 0$$

OFF-SHELL

Step 3.

Feynman gauge $\xi = 1$

Integrate out field **b**

$$\mathcal{L} = \frac{1}{2} \phi^a \square \phi^a + \bar{c} \square c$$

$$\mathbf{Z} = \frac{D_0^2}{D_1}$$

BRST

$$s\phi^a = \partial^a c$$

$$sc = 0$$

$$s\bar{c} = \mathbf{L}$$

anti-BRST

$$\bar{s}\phi^a = \partial^a \bar{c}$$

$$\bar{s}c = -\mathbf{L}$$

$$\bar{s}\bar{c} = 0$$

$$s^2 = 0 \quad \bar{s}^2 = 0$$

ON-SHELL for c and \bar{c}

spin-2 in $R^{3,1}$ (Weyl gravity)

$$\begin{array}{cc} \phi_{-1}^{ab} & \phi_1^{ab} \\ & \phi_0^a \end{array}$$

$$\Delta(\phi_{k'}) = \frac{d-2}{2} + k'$$

$$\begin{aligned}
\mathcal{L} &= \frac{1}{2}\phi_{-1}^{ab}\square\phi_1^{ab} - \frac{1}{4}\phi_{-1}^{aa}\square\phi_1^{bb} + \frac{1}{2}\phi_0^a\square\phi_0^a \\
&+ L_{-1}^a L_1^a + \frac{1}{2}L_0 L_0 \\
&- \frac{1}{4}\phi_1^{ab}\phi_1^{ab} + \frac{1}{8}\phi_1^{aa}\phi_1^{bb}
\end{aligned}$$

$$L_{-1}^a = \partial^b \phi_{-1}^{ab} - \frac{1}{2}\partial^a \phi_{-1}^{bb} + \phi_0^a$$

$$L_1^a = \partial^b \phi_1^{ab} - \frac{1}{2}\partial^a \phi_1^{bb}$$

$$L_0 = \partial^a \phi_0^a + \frac{1}{2}\phi_1^{bb}$$

$$\delta\phi_{-1}^{ab} = \partial^a \xi_{-1}^b + \partial^b \xi_{-1}^a + \eta^{ab} \xi_0$$

$$\delta\phi_1^{ab} = \partial^a \xi_1^b + \partial^b \xi_1^a$$

$$\delta\phi_0^a = \partial^a \xi_0 - \xi_1^a$$

Faddeev-Popov fields

$$\bar{c}_{-1}^a \quad c_{-1}^a, \quad \bar{c}_1^a \quad c_1^a \quad \bar{c}_0 \quad c_0$$

Nakanishi-Lautrup field

$$b_{-1}^a \quad b_1^a, \quad b_0$$

$$\mathcal{L}^{BRST} = \mathcal{L} + \mathcal{L}_{\text{g.fix}}$$

$$\begin{aligned} \mathcal{L}_{\text{g.fix}} = & -b_1^a L_{-1}^a - b_{-1}^a L_1^a - b_0 L_0 \\ & + \xi b_1^a b_{-1}^a + \frac{1}{2} \xi b_0^2 \\ & + \bar{c}_1^a \square c_{-1}^a + \bar{c}_{-1}^a \square c_1^a + \bar{c}_0 \square c_0 \\ & - \bar{c}_1^a c_1^a \end{aligned}$$

Use Feynman gauge $\xi = 1$

Integrate out Nakanishi-Lautrup fields

$$\begin{aligned}
\mathcal{L}^{BRST} &= \frac{1}{2}\phi_{-1}^{ab}\square\phi_1^{ab} - \frac{1}{4}\phi_{-1}^{aa}\square\phi_1^{bb} + \frac{1}{2}\phi_0^a\square\phi_0^a \\
&- \frac{1}{4}\phi_1^{ab}\phi_1^{ab} + \frac{1}{8}\phi_1^{aa}\phi_1^{bb} \\
&+ \bar{c}_1^a\square c_{-1}^a + \bar{c}_{-1}^a\square c_1^a + \bar{c}_0\square c_0 \\
&- \bar{c}_1^a c_1^a
\end{aligned}$$

Solution for **auxiliary fields**

$$\phi_1^{ab} = \square \phi_{-1}^{ab}$$

$$c_1^a = \square c_{-1}^a$$

$$\bar{c}_1^a = \square \bar{c}_{-1}^a$$

$$\mathcal{L} = \frac{1}{2} \phi_{-1}^{ab} \square^2 \phi_{-1}^{ab} - \frac{1}{4} \phi_{-1}^{aa} \square^2 \phi_{-1}^{bb} + \frac{1}{2} \phi_0^a \square \phi_0^a$$

$$+ \bar{c}_{-1}^a \square^2 c_{-1}^a + \bar{c}_0 \square c_0$$

$$Z = \frac{D_1^4 D_0^2}{D_2^2 D_0^2 D_1}$$

$$= \frac{D_1^3}{D_2^2}$$

$$D_n = \sqrt{\det(-\square)}$$

Arbitrary spin- s conformal field

double traceless fields

$$\phi_{\lambda, k'}^{a_1 \dots a_{s'}}$$

$$s' = 0, 1, \dots, s$$

$$\lambda \in [s - s']_2$$

$$k' \in [\kappa - 1 + \lambda]_2$$

$$p \in [q]_2 \iff p = -q, -q + 2, \dots, q - 2, q$$

Lagrangian

$$\mathcal{L} = \sum_{\lambda, k', s'} \mathcal{L}_{\lambda, k'}^{s'}$$

$$\begin{aligned} \mathcal{L}_{\lambda, k'}^{s'} &= \phi^{s'} \square \phi^{s'} + \phi^{s'} \phi^{s'} \\ &+ \mathbf{L}^{s'} \mathbf{L}^{s'} \end{aligned}$$

$$\begin{aligned} \mathbf{L}^{s'} &= \partial^{\mathbf{a}} \phi^{\mathbf{a} \mathbf{a}_1 \dots \mathbf{a}_{s'}} - \frac{1}{2} \partial^{\mathbf{a}_1} \phi^{\mathbf{a} \mathbf{a}_2 \dots \mathbf{a}_{s'}} \\ &+ \phi^{\mathbf{a}_1 \dots \mathbf{a}_{s'}} \\ &+ \eta(\mathbf{a}_1 \mathbf{a}_2 \phi^{\mathbf{a}_3 \dots \mathbf{a}_{s'}}) \end{aligned}$$

$$\begin{aligned}\delta\phi^{a_1\dots a_{s'}} &= \partial(a_1\xi^{a_2\dots a_{s'}}) \\ &+ \xi^{a_1\dots a_{s'}} \\ &+ \eta^{(a_1a_2\xi^{a_3\dots a_{s'}})}\end{aligned}$$

Apply Faddeev-Popov procedure

Faddeev-Popov fields

$$\bar{c}^{a_1 \dots a_{s'}} \quad c^{a_1 \dots a_{s'}}$$

Nakanishi-Lautrup fields

$$b^{a_1 \dots a_{s'}}$$

$$\mathcal{L}^{BRST} = \mathcal{L} + \mathcal{L}_{\text{g.fix}}$$

$$\mathcal{L}_{\text{g.fix}} = \sum_{s'} \mathcal{L}_{\text{g.fix}}^{s'}$$

$$\begin{aligned} \mathcal{L}_{\text{g.fix}}^{s'} &= b^{s'} \mathbf{L}^s + \frac{1}{2} \xi \mathbf{b}^{s'} \mathbf{b}^{s'} \\ &+ \bar{\mathbf{c}}^{s'} \square \mathbf{c}^{s'} \\ &+ \bar{\mathbf{c}}^{s'} \mathbf{c}^{s'} \end{aligned}$$