

Based on 1603.05387

Collaboration with [Jin-beom BAE](#) and [Shailesh LAL](#)

One Loop Test of Free Adjoint Model Holography

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(Free) Vector Model Duality

Sezgin, Sundell, Klebanov, Polyakov

O(N)/U(N)
Vector Models

Single trace operators

$$T^{\mu\nu}$$

$$J^\mu$$

$$J^{\mu_1 \cdots \mu_s}$$

~~$$\mathcal{O}_{\Delta}^{\mu_1 \cdots \mu_s}$$~~

Vasiliev Theories

Fields (single particles)

$$G_{MN}$$

$$A_M$$

$$\varphi_{M_1 \cdots M_s}$$

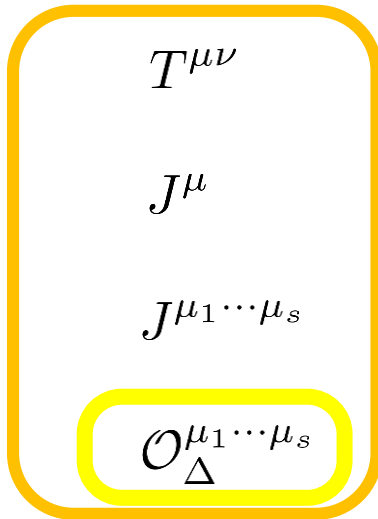
~~$$\phi_{M_1 \cdots M_s}$$~~

Free Adjoint Scalar Model Duality

Free Adjoint Scalar

$$S_{\text{CFT}}[\phi] = \int d^d x \text{Tr} \left[\phi^\dagger \square \phi \right]$$

Single trace operators



$$\text{Tr} \left[(\partial_{\mu_1} \partial_{\mu_2} \dots \partial_{\mu_l} \phi) (\partial_{\nu_1} \partial_{\nu_2} \dots \partial_{\nu_m} \phi) \dots (\partial_{\rho_1} \partial_{\rho_2} \dots \partial_{\rho_n} \phi) \right]$$

HS Gauge Theories
+ HS Matters

Fields (single particles)

G_{MN}

What Can We Do

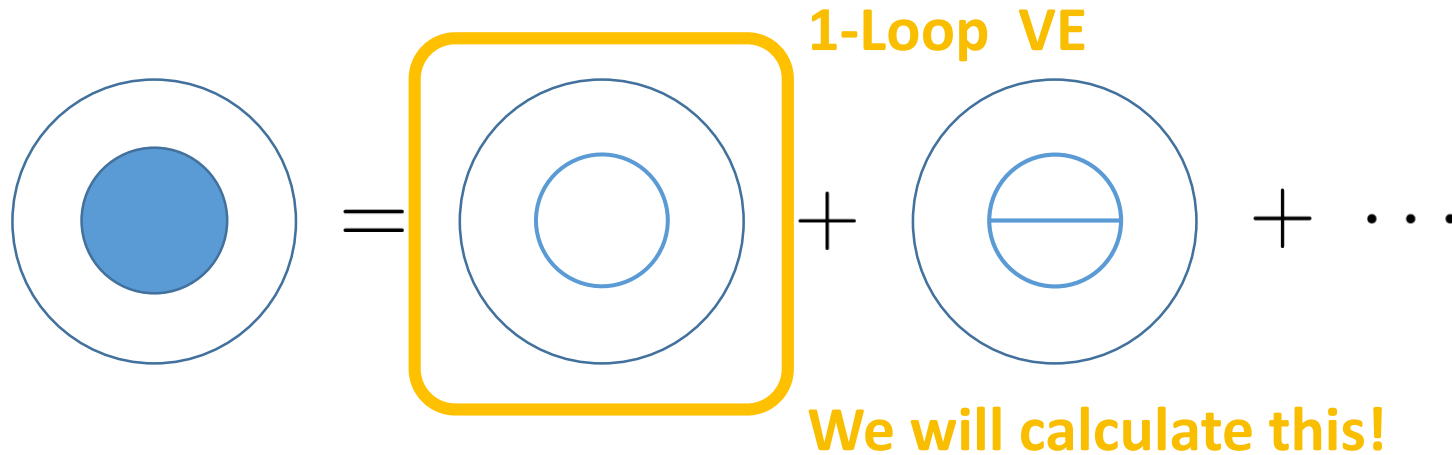
A_M

About This Theory?

$\varphi_{M_1 \dots M_s}$

$\phi_{M_1 \dots M_s}$

No External Leg: **AdS Vacuum (or Free) Energy**



- No **interaction** involved
- We only need the **spectrum** of AdS theory
- Expect to vanish

Vector Model Case

- Calculated by [Giombi, Klebanov, Safidi; Tseytlin \[2013-...\]](#)
- VE vanishes for non-min HS, but not for min HS (shift $N \rightarrow N-1$)

$$\begin{aligned}
 \text{Bubble} &= \text{Bubble}(\phi) + \text{Bubble}(A_M) + \text{Bubble}(G_{MN}) + \text{Bubble}(\varphi_{MNL}) + \dots \\
 &= \sum_{\Delta, \ell} N_{\Delta, \ell} \text{Bubble}(\phi_{\Delta, \ell})
 \end{aligned}$$

We need to know

- Multiplicity $N_{\Delta, s}$

- Vacuum (or Free) Energy of AdS Field: $\Gamma_{\Delta, \ell}^{(1)}(z) = \text{Bubble}(\phi_{\Delta, \ell})_{\Lambda_{UV}}$ $\frac{1}{z} = \log(\Lambda_{UV} R)$

❖ Zeta Function

$$\Gamma_{\Delta, \ell}^{(1)}(\Lambda) = -\frac{1}{2} \Gamma(z) \zeta_{\Delta, \ell}(z), \quad \zeta_{\Delta, \ell}(z) = \int_0^\infty \frac{dt}{t} \frac{t^z}{\Gamma(z)} K_{\Delta, \ell}(t)$$

Zeta Function (AdS with Sphere bd)

- UV divergence $\Gamma_{\Delta,\ell}^{(1)\text{div}} = -\frac{\zeta_{\Delta,\ell}(0)}{2} \log \frac{\Lambda_{\text{UV}}}{\mu}$
- Renormalized VE $\Gamma_{\Delta,\ell}^{(1)\text{ren}} = -\frac{\zeta_{\Delta,\ell}(0)}{2} \log(\mu R) - \frac{\zeta'_{\Delta,\ell}(0)}{2}$
- Integral Representation for $\zeta_{\Delta,\ell}(z)$ Camporesi, Higuchi; Beccaria, Tseytlin

Vector Model Case

- Simple Multiplicity (Flato-Fronsdal)

- A** {
- UV divergent VE $\sum_s \zeta_{s+d-2,s}(0) = 0 + \infty$
 - Renormalized VE $\sum_s \zeta'_{s+d-2,s}(0) = 0 + \infty$
- Additional Regularization
- B** {
- Total Zeta Function $\zeta_{\text{AdS}}(z) = \sum_s \zeta_{s+d-2,s}(z) = \mathcal{O}(z^2)$

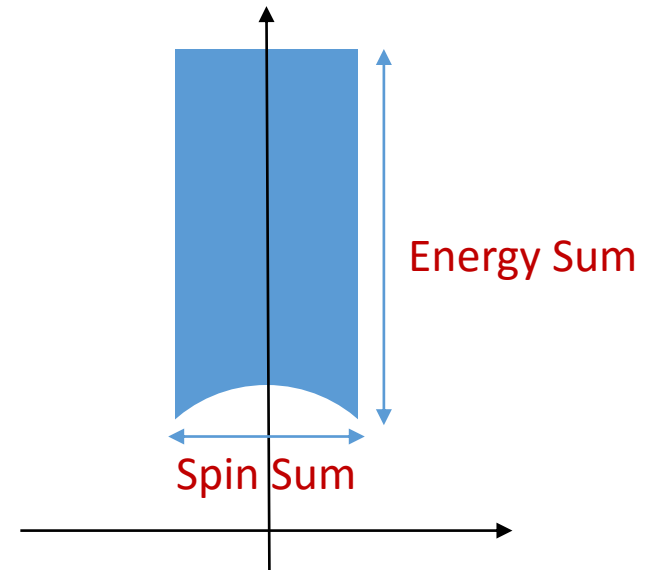
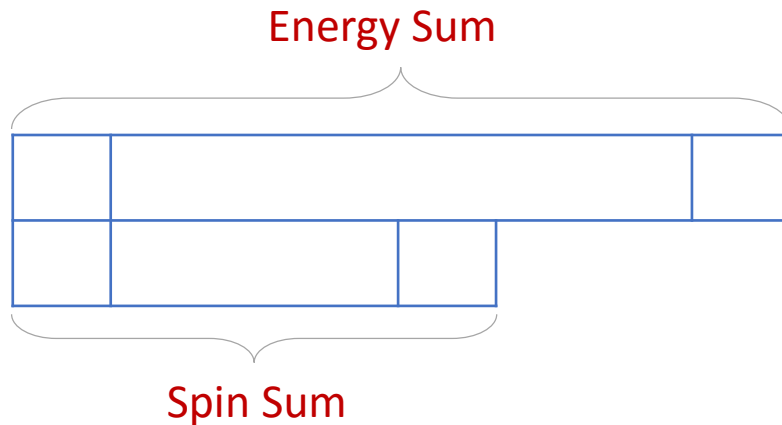
Summation Order of Energy vs Spin

A

- 1) Energy Sum : **Divergent** → **Regularization**
- 2) Spin Sum : **Divergent** → **Regularization**

B

- 1) Spin Sum : **Convergent**
- 2) Energy Sum : **Convergent**



Regularization controlled by HS Symmetry?

VE vanishes for each HS multiplet?



I. Single Particle **Multiplicities**
of Each **Massive HS Multiplets**

II. **Vacuum Energy**
of Each **Massive HS Multiplets**





I. Single Particle **Multiplicities**
of Each **Massive HS Multiplets**

(Singleton)¹ 'Fund. Rep.' of HS Sym

$$\partial_{\mu_1} \partial_{\mu_2} \cdots \partial_{\mu_l} \phi \quad \longleftrightarrow \quad \text{Rac} = \mathcal{D}(\tfrac{1}{2}, 0)$$

(Singleton)² Massless HS Multiplet

$$(\partial_{\mu_1} \partial_{\mu_2} \cdots \partial_{\mu_l} \phi) (\partial_{\nu_1} \partial_{\nu_2} \cdots \partial_{\nu_m} \phi) \quad \longleftrightarrow \quad \text{Rac} \otimes \text{Rac}$$

(Singleton)ⁿ Massive HS Multiplet

$$(\partial_{\mu_1} \partial_{\mu_2} \cdots \partial_{\mu_l} \phi) (\partial_{\nu_1} \partial_{\nu_2} \cdots \partial_{\nu_m} \phi) \cdots (\partial_{\rho_1} \partial_{\rho_2} \cdots \partial_{\rho_n} \phi) \quad \longleftrightarrow \quad \text{Rac}^{\otimes n}$$

Flato-Fronsdal Theorem (F^2 Thm.)

$$\text{Rac}^{\otimes 2} = \bigoplus_{s=0}^{\infty} \mathcal{D}(s+1, s), \quad \text{Di}^{\otimes 2} = \mathcal{D}(2, 0) \oplus \bigoplus_{s=1}^{\infty} \mathcal{D}(s+1, s)$$

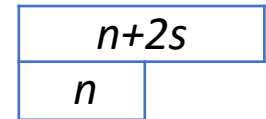
massless spin s
in AdS₄

Higher Power Extension (F^n Thm.)

- Oscillator Method [Konstein, Vasiliev](#)

$$(\text{Di} \oplus \text{Rac})^{\otimes p} = \bigoplus_{n=0}^{\infty} \bigoplus_{2s=0}^{\infty} \dim \left(\pi_{(n+2s, n)}^{O(p)} \right) \mathcal{D} \left(s + n + \frac{p}{2}, s \right)$$

$$\text{Di}^{\otimes q} \otimes \text{Rac}^{\otimes p-q} = \bigoplus_{n=0}^{\infty} \bigoplus_{2s=0}^{\infty} N_{(n+2s, n)}^{[q, p-q]} \mathcal{D} \left(s + n + \frac{p}{2}, s \right)$$



- Character Method [Barabanschikov, Grant, Huang, Raju; Newton, Spradlin](#)

- Concrete Cases

Order Three (F^3) or '2nd Regge Trajectory'

$$\text{Rac}^{\otimes 3} = \bigoplus_{s=0}^{\infty} (s+1) \left[\mathcal{D}\left(s + \frac{3}{2}, s\right) \oplus \mathcal{D}\left(s + \frac{7}{2}, s+1\right) \right]$$

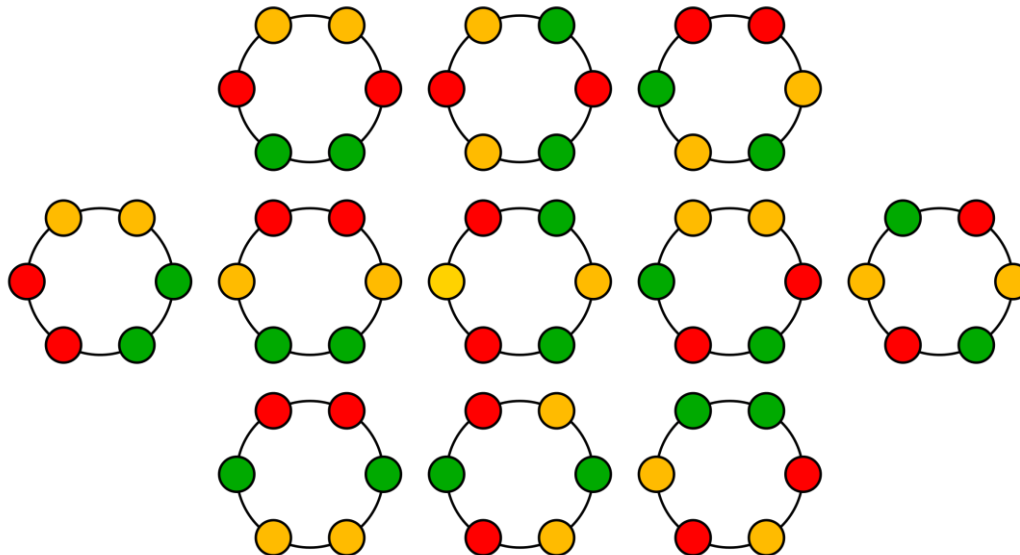
Order Four (F^4) or '3rd Regge Trajectory'

$$\begin{aligned} \text{Rac}^{\otimes 4} = & \bigoplus_{s=0}^{\infty} \frac{(1+s)(2+s)}{2} \mathcal{D}(s+2, s) \\ & \oplus \bigoplus_{s=0}^{\infty} \bigoplus_{n=1}^{\infty} \frac{(2n+2s+1)(2s+1) + 3(-1)^n}{4} \mathcal{D}(s+n+2, s) \end{aligned}$$

Single Trace Operators

$$\text{Tr} \left[(\partial_{\mu_1} \partial_{\mu_2} \cdots \partial_{\mu_l} \phi) (\partial_{\nu_1} \partial_{\nu_2} \cdots \partial_{\nu_m} \phi) \cdots (\partial_{\rho_1} \partial_{\rho_2} \cdots \partial_{\rho_n} \phi) \right]$$

Counting Number of Necklaces



Polya Enumeration Thm. [Polyakov; Bianchi, Morales, Samtleben, Beisert; Spradlin, Volovich](#)

SU(N) Trace : Cyclic Tensor Product

$$\chi_{\text{cyc}^n}(g) = \frac{1}{n} \sum_{k|n} \varphi(k) \left(\chi_V(g^k) \right)^{\frac{n}{k}}$$

Euler totient fn:

of relative prime of k in {1,...,k}

$$T_{\text{cyc}}^{(2)}(\text{Rac}) = \bigoplus_{n=0}^{\infty} \mathcal{D}(2n+1, 2n)$$

$$T_{\text{cyc}}^{(3)}(\text{Rac}) = \bigoplus_{s=0}^{\infty} \left(s+1+2\left[-\frac{s}{3}\right] \right) \left[\mathcal{D}\left(s+\frac{3}{2}, s\right) \oplus \mathcal{D}\left(s+\frac{7}{2}, s+1\right) \right]$$

Order four formula does not fit in a single page



II. Vacuum Energy of Each Massive HS Multiplets

A

- UV divergent VE

$$\sum_{\Delta,s} N_{\Delta,s} \zeta_{\Delta,s}(0)$$

- Renormalized VE

$$\sum_{\Delta,s} N_{\Delta,s} \zeta'_{\Delta,s}(0)$$

B

- Total Zeta Function

$$\zeta_{\text{AdS}}(z) = \sum_{\Delta,s} N_{\Delta,s} \zeta_{\Delta,s}(z)$$

A Toy Model : Rac³

$$\text{Rac}^{\otimes 3} = \bigoplus_{s=0}^{\infty} (s+1) \left[\mathcal{D}(s + \frac{3}{2}, s) \oplus \mathcal{D}(s + \frac{7}{2}, s+1) \right]$$

$$\zeta_{\Delta,s}(0) = \frac{2s+1}{24} \left[\left(\Delta - \frac{3}{2} \right)^4 - \left(\frac{2s+1}{2} \right)^2 \left(2 \left(\Delta - \frac{3}{2} \right)^2 + \frac{1}{6} \right) - \frac{7}{240} \right]$$

$$\zeta'_{\Delta,s}(0) = \frac{2s+1}{3} \left[\frac{1}{8} \left(\Delta - \frac{3}{2} \right)^4 + \frac{1}{48} \left(\Delta - \frac{3}{2} \right)^2 + c_3 + \left(\frac{2s+1}{2} \right)^2 c_1 + \int_0^{\Delta - \frac{3}{2}} dx \left(\left(\frac{2s+1}{2} \right)^2 - x^2 \right) x \psi \left(x + \frac{1}{2} \right) \right].$$

$$c_n = \int_0^{\infty} du \frac{2u^n \ln u}{e^{2\pi u} + 1}, \quad \psi(x) = \int_0^{\infty} dt \left(\frac{e^{-t}}{t} - \frac{e^{-xt}}{1 - e^{-t}} \right)$$



A Toy Model : Rac³

UV divergent part of Vacuum Energy

$$\begin{aligned}\zeta_{\text{Rac}^{\otimes 3}}(0) &= \lim_{\alpha \rightarrow 0} \sum_{s=0}^{\infty} \left(\Delta - \frac{3}{2}\right)^{-\alpha} \zeta_{\Delta, s}(0) \\ &= \lim_{\alpha \rightarrow 0} \sum_{s=0}^{\infty} s^{-\alpha} \zeta_{s+\frac{3}{2}, s}(0) + \lim_{\alpha \rightarrow 0} \sum_{s=0}^{\infty} (s+1)^{-\alpha} \zeta_{s+\frac{5}{2}, s}(0)\end{aligned}$$



•
•
•


$$= -\frac{7}{13824} + \frac{7}{13824} = 0$$



A Toy Model : Rac³

Renormalized Vacuum Energy

$$\zeta'_{\text{Rac}^{\otimes 3}}(0) = \lim_{\alpha \rightarrow 0} \sum_{s=0}^{\infty} \int_0^{\infty} dt \left[s^{-\alpha} (s+1) \int_0^s dx + (s+1)^{-\alpha} s \int_0^{s+1} dx \right] \times$$



$$\times \frac{2}{3} \left(\frac{2s+1}{2} \right) \left[\left(\frac{2s+1}{2} \right)^2 x - x^3 \right] \left(\frac{e^{-t}}{t} - \frac{e^{-(x+\frac{1}{2})t}}{1-e^{-t}} \right)$$

⋮

$$= -\frac{\ln 2}{128} - \frac{19 \zeta(3)}{3840 \pi^2} + \frac{25 \zeta(5)}{256 \pi^4} - \frac{63 \zeta(7)}{512 \pi^6}$$



Rac⁴

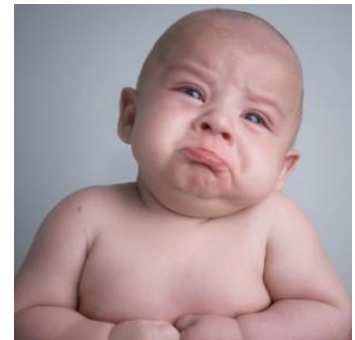
UV divergent part of Vacuum Energy

$$\begin{aligned} \zeta_{\text{Rac}^{\otimes 4}}(0) &= \lim_{\alpha \rightarrow 0} \sum_{s=0}^{\infty} \sum_{n=1}^{\infty} \left(s + n + \frac{1}{2}\right)^{-\alpha} \frac{2n(2s+1) + 3(-1)^n + (2s+1)^2}{4} \zeta_{s+n+2,s}(0) \\ &\quad + \lim_{\alpha \rightarrow 0} \sum_{s=0}^{\infty} \left(s + \frac{1}{2}\right)^{-\alpha} \frac{(s+1)(s+2)}{2} \zeta_{s+2,s}(0) \\ &= \dots = -\frac{7963}{232243200} + \frac{7963}{232243200} = 0 \end{aligned}$$



Renormalized Vacuum Energy

$$\begin{aligned} \zeta'_{\text{Rac}^{\otimes 4}}(0) &= \dots \\ &= -\frac{29 \zeta(3)}{2520 \pi^2} - \frac{\zeta(5)}{60 \pi^4} + \frac{\zeta(7)}{8 \pi^6} - \frac{\zeta(9)}{8 \pi^8} \end{aligned}$$



Cyclic Tensor Products ?

Hard to control the regularization scheme

No Way to Proceed ?



Rac^n

Character

Decomposition

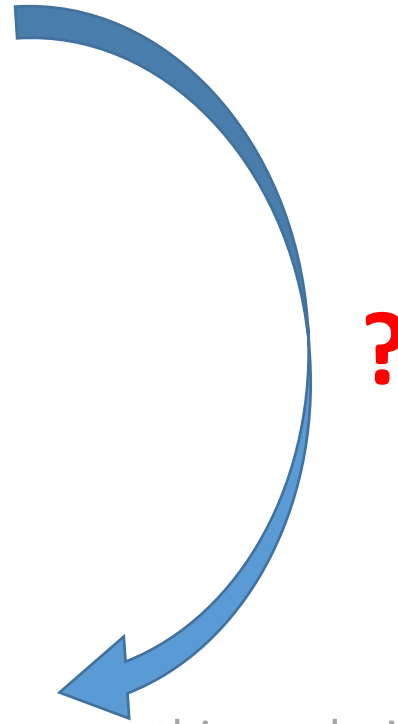


Multiplicities

Resummation



Zeta Function



this works in Thermal AdS
Giombi, Klebanov, Tseytlin

New Method to Calculate Zeta Function



We introduce

Character Integral Representation for Zeta Function

- For any spectra \mathcal{H}
- Identify the corresponding character $\chi_{\mathcal{H}}(\beta, \alpha)$
- Total zeta function of \mathcal{H} given by



$$\tilde{\zeta}_{\mathcal{H}}(z) = \int_0^{\infty} d\beta \frac{\beta^{2z-1}}{\Gamma(2z)} \frac{\cosh \frac{\beta}{2}}{\sinh \frac{\beta}{2}} \left(1 + \sinh^2 \frac{\beta}{2} \partial_{\alpha}^2 \right) \chi_{\mathcal{H}}(\beta, \alpha) \Big|_{\alpha=0}$$

Great simplification !

since we usually identify the spectra using character

Revisit (test of formula)

Vector Model

$$\tilde{\zeta}_{\text{non-min}}(z) = 0 \quad \tilde{\zeta}_{\text{min}}(z) = 4^{-z} \tilde{\zeta}_{\text{Rac}}(z) = \tilde{\zeta}_{\text{Rac}}(z) + \mathcal{O}(z^2)$$

integrand itself vanishes

$$\tilde{\zeta}_{\text{Rac}}(z) = \zeta\left(2z - 2, \frac{1}{2}\right) + \frac{1}{4} \zeta\left(2z, \frac{1}{2}\right) = - \left[\frac{\ln 2}{4} - \frac{3 \zeta(3)}{8 \pi^2} \right] z + \mathcal{O}(z^2)$$

Coincide with bd 1-loop calculation
more discussion: [Becarria](#), [Tseytlin](#)

Rac³ and Rac⁴

$$\begin{aligned} \tilde{\zeta}_{\text{Rac}^{\otimes 3}}(z) &= \frac{\zeta\left(2z, \frac{3}{2}\right)}{128} - \frac{19 \zeta\left(2z - 2, \frac{3}{2}\right)}{1440} - \frac{5 \zeta\left(2z - 4, \frac{3}{2}\right)}{72} - \frac{\zeta\left(2z - 6, \frac{3}{2}\right)}{90} \\ &= z \left(-\frac{\ln 2}{128} - \frac{19 \zeta(3)}{3840 \pi^2} + \frac{25 \zeta(5)}{256 \pi^4} - \frac{63 \zeta(7)}{512 \pi^6} \right) + \mathcal{O}(z^2), \end{aligned}$$

$$\begin{aligned} \tilde{\zeta}_{\text{Rac}^{\otimes 4}}(z) &= \frac{29 \zeta(2z - 2)}{1260} - \frac{\zeta(2z - 4)}{90} - \frac{\zeta(2z - 6)}{90} - \frac{\zeta(2z - 8)}{1260} \\ &= z \left(-\frac{29 \zeta(3)}{2520 \pi^2} - \frac{\zeta(5)}{60 \pi^4} + \frac{\zeta(7)}{8 \pi^6} - \frac{\zeta(9)}{8 \pi^8} \right) + \mathcal{O}(z^2). \end{aligned}$$

Coincide with the previous calculations **(A)**

Free SU(N) Adjoint Scalar : '2nd Regge Trajectory'

$$\begin{aligned} \tilde{\zeta}_{\text{cyc}^3}(z) = & \frac{\zeta\left(2z, \frac{3}{2}\right)}{1152} - \frac{19\zeta\left(2z-2, \frac{3}{2}\right)}{17280} - \frac{5\zeta\left(2z-4, \frac{3}{2}\right)}{216} - \frac{\zeta\left(2z-6, \frac{3}{2}\right)}{270} \\ & + 2^{2z} 3^{-2z-1} [\zeta(2z) + \zeta(2z-2)] - 3^{-2z-1} [\zeta(2z) + 4\zeta(2z-2)] \\ & + 3^{-2z-3} \frac{57}{32} [\zeta\left(2z, \frac{1}{6}\right) + \zeta\left(2z, \frac{5}{6}\right) - 8\zeta\left(2z-1, \frac{1}{6}\right) + 8\zeta\left(2z-1, \frac{5}{6}\right)] \\ & + 3^{-2z-1} [\zeta\left(2z-2, \frac{1}{6}\right) + \zeta\left(2z-2, \frac{5}{6}\right) - 3\zeta\left(2z-3, \frac{1}{6}\right) + 3\zeta\left(2z-3, \frac{5}{6}\right)] \end{aligned}$$



UV divergent part of Vacuum Energy

$$\tilde{\zeta}_{\text{cyc}^3}(0) = 0 \quad \text{😊}$$

Renormalized Vacuum Energy

$$\begin{aligned} \tilde{\zeta}_{\text{cyc}^3}'(0) &= -\frac{43 \ln 2}{128} + \frac{1487 \zeta(3)}{3840 \pi^2} + \frac{25 \zeta(5)}{768 \pi^4} - \frac{21 \zeta(7)}{512 \pi^6} + \frac{4 \pi}{27\sqrt{3}} - \frac{19 \psi^{(1)}\left(\frac{1}{3}\right)}{72\sqrt{3} \pi} + \frac{\psi^{(3)}\left(\frac{1}{3}\right)}{96\sqrt{3} \pi^3} \\ &\simeq -0.311588, \quad \text{😞} \end{aligned}$$

'3rd Regge Trajectory'

UV divergent part of Vacuum Energy

$$\tilde{\zeta}_{\text{cyc}^4}(0) = 0 \quad \text{😊}$$

Renormalized Vacuum Energy

$$\begin{aligned} \tilde{\zeta}_{\text{cyc}^4}'(0) &= -\frac{25 \ln 2}{64} - \frac{4573 \zeta(3)}{20160 \pi^2} + \frac{457 \zeta(5)}{1920 \pi^4} + \frac{\zeta(7)}{32 \pi^6} - \frac{\zeta(9)}{32 \pi^8} + \frac{\pi}{32} - \frac{5 \psi^{(1)}\left(\frac{1}{4}\right)}{96 \pi} + \frac{\psi^{(3)}\left(\frac{1}{4}\right)}{384 \pi^3} \\ &\simeq -0.353518. \end{aligned} \quad \text{😞}$$

Conjecture of Free **Adjoint** Model Holography:

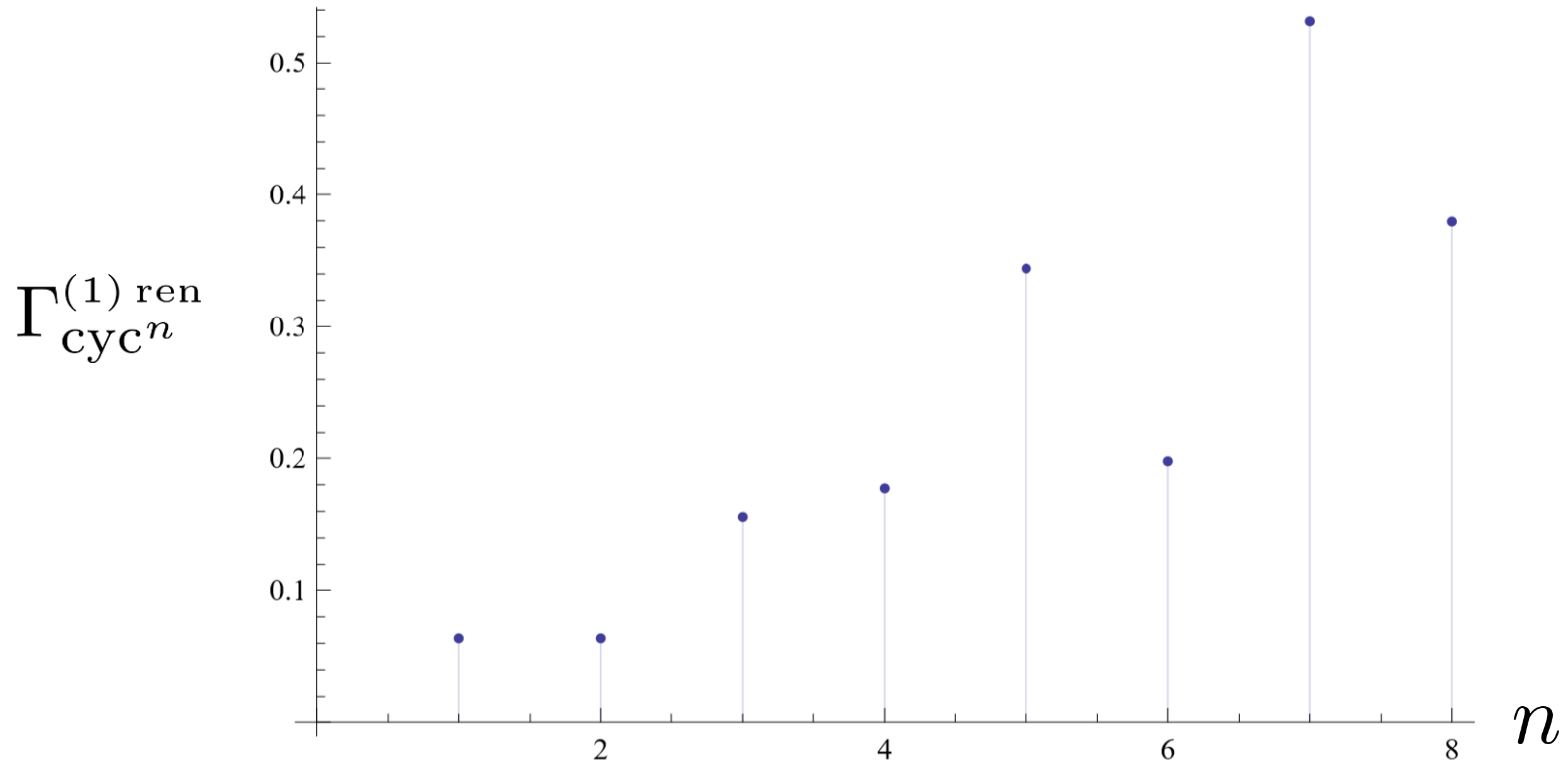


Maybe, we need to sum over *all trajectories*

Q

- Why each massive **HS multiplet** do not show any special property?
- No particular role of **HS sym**?
- What about **Multiplet HS sym**?

VE of First Few Trajectories



pattern not so clear

AdS₅/CFT₄

Character Integral Representation for Zeta Function

$$\zeta_{\mathcal{H}}(z) := \zeta_{\mathcal{H}|2}(z) + \zeta_{\mathcal{H}|1}(z) + \zeta_{\mathcal{H}|0}(z)$$

$$\frac{\Gamma(z) \zeta_{\mathcal{H}|n}(z)}{\log R} = \int_0^\infty d\beta \frac{\left(\frac{\beta}{2}\right)^{2(z-1-n)}}{\Gamma(z-n)} f_{\mathcal{H}|n}(\beta)$$

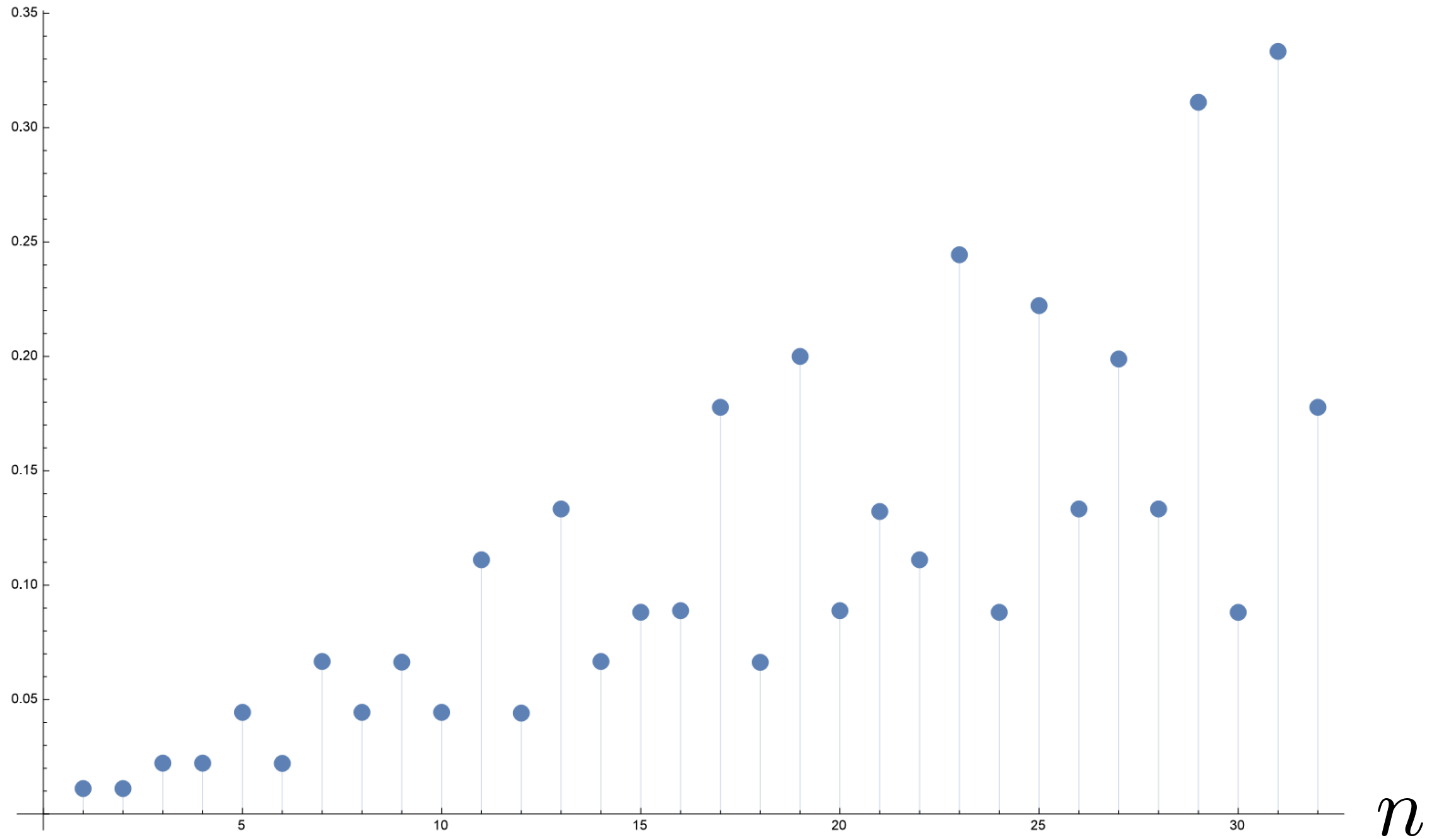
$$f_{\mathcal{H}|2}(\beta) = \frac{\sinh^4 \frac{\beta}{2}}{2} \chi_{\mathcal{H}}(\beta, 0, 0)$$

$$f_{\mathcal{H}|1}(\beta) = \sinh^2 \frac{\beta}{2} \left[\frac{\sinh^2 \frac{\beta}{2}}{3} - 1 - \sinh^2 \frac{\beta}{2} (\partial_{\alpha_1}^2 + \partial_{\alpha_2}^2) \right] \chi_{\mathcal{H}}(\beta, \alpha_1, \alpha_2) \Big|_{\alpha_i=0}$$

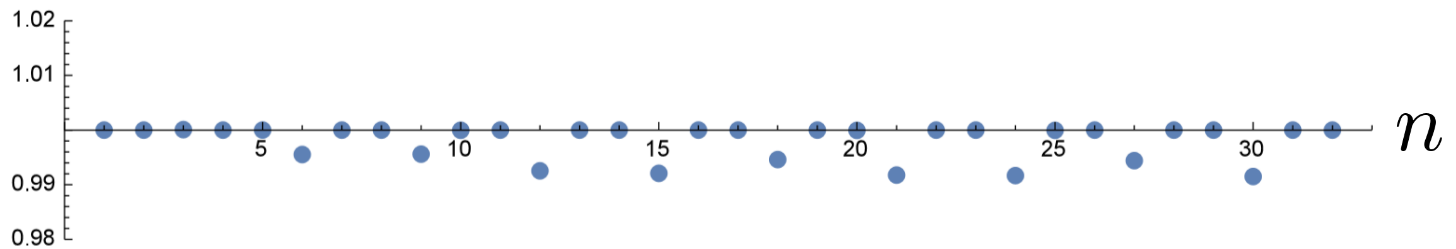
$$f_{\mathcal{H}|0}(\beta) = \left[1 + \frac{\sinh^2 \frac{\beta}{2} (3 - \sinh^2 \frac{\beta}{2})}{3} (\partial_{\alpha_1}^2 + \partial_{\alpha_2}^2) - \frac{\sinh^4 \frac{\beta}{2}}{3} (\partial_{\alpha_1}^4 - 12 \partial_{\alpha_1}^2 \partial_{\alpha_2}^2 + \partial_{\alpha_2}^4) \right] \chi_{\mathcal{H}}(\beta, \alpha_1, \alpha_2) \Big|_{\alpha_i=0}$$

VE of First Few Trajectories

$$\frac{\Gamma_{\text{cyc}^n}^{(1) \text{ ren}}}{\log R}$$



$$\frac{\Gamma_{\text{cyc}^n}^{(1) \text{ ren}}}{\varphi(n) \Gamma_{\text{Rac}}^{(1) \text{ ren}}}$$



Certain pattern. But not yet clear how to deal with it

Slicing Vacuum Energy Differently

Total VE in Different Slicings

Sum over Trajectories

$$\chi_{\text{adj}} = \sum_{n=2}^{\infty} \chi_{\text{cyc}^n}$$

Sum over 'log slice's

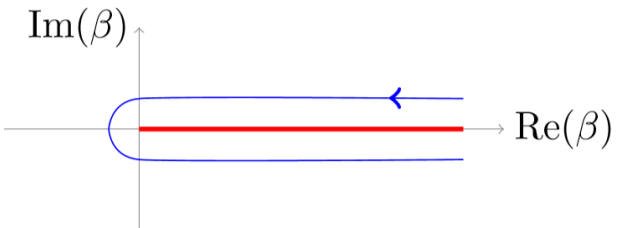
$$\chi_{\text{adj}} = -\chi_{\text{Rac}} + \sum_{k=1}^{\infty} \frac{\varphi(k)}{k} \chi_{\text{log},k}$$

$$\chi_{\text{log},k}(\beta, \alpha_1, \alpha_2) = -\log[1 - \chi_{\text{Rac}}(k\beta, k\alpha_1, k\alpha_2)]$$

A Subtlety

$$\frac{\Gamma(z) \zeta_{\mathcal{H}|n}(z)}{\log R} = \int_0^\infty d\beta \frac{\left(\frac{\beta}{2}\right)^{2(z-1-n)}}{\Gamma(z-n)} f_{\mathcal{H}|n}(\beta)$$

- 1 Extract β^{2n+1} coefficients from $f_{\mathcal{H}|n}(\beta)$
- 2 Integrate with analytic continuation on z
- 3 Contour integral

$$\frac{i}{2 \sin(2\pi z)} \oint_C d\beta \frac{\left(-\frac{\beta}{2}\right)^{2(z-1-n)}}{\Gamma(z-n)} f_{\mathcal{H}|n}(\beta)$$


They are **equivalent** for **single ptcl**
and for each **trajectory**, but **not** for each 'log slice'

$$\chi_{\log,k}(\beta, \alpha_1, \alpha_2) = -\log[1 - \chi_{\text{Rac}}(k\beta, k\alpha_1, k\alpha_2)]$$

1

$$f_{\log,k|2}(\beta) = \frac{-3 \log \beta + \log(-\frac{2}{k^3})}{32} \beta^4 + \frac{-3 \log \beta + \log(-\frac{2}{k^3})}{192} \beta^6 - \frac{k^3}{64} \beta^7 + \mathcal{O}(\beta^8)$$

No β^5

$$f_{\log,k|1}(\beta) = \frac{3 \log \beta - \log(-\frac{2}{k^3}) + 1}{8} \beta^2 - \frac{k^2 - 2}{96} \beta^4 + \frac{k^3}{8} \beta^5 + \mathcal{O}(\beta^6)$$

No β^3

$$f_{\log,k|0}(\beta) = -3 \log \beta + \log(-\frac{2}{k^3}) - \frac{3}{2} + \frac{k^2 - 1}{12} \beta^2 - \frac{3k^2}{2} \beta^3 + \mathcal{O}(\beta^4)$$

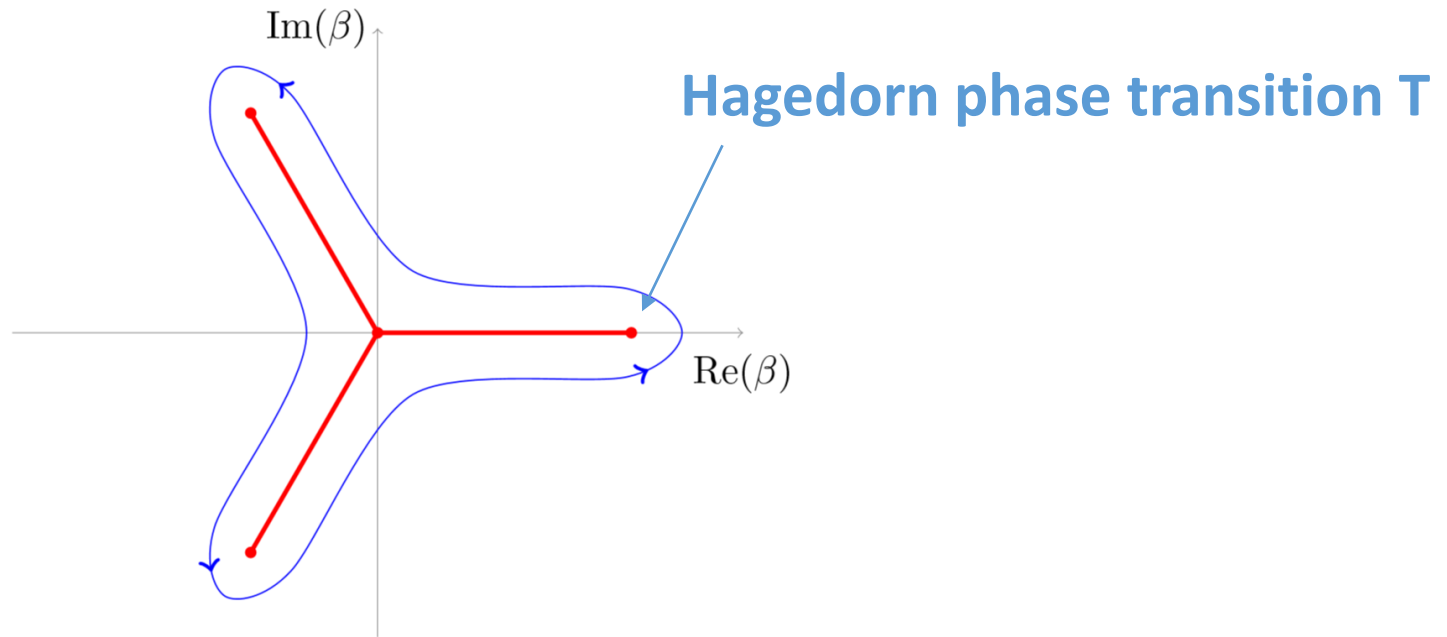
No β

2 can be done using an additional regularization

Vacuum Energy vanishes for each k 'log slice' !

$$\chi_{\log,k}(\beta, \alpha_1, \alpha_2) = -\log[1 - \chi_{\text{Rac}}(k\beta, k\alpha_1, k\alpha_2)]$$

3



But, β is integration parameter here !

Any interpretation ?

OUTLOOK

- ❖ $SO(N)$ & $Sp(N)$ Adjoint, Bivector Models
- ❖ Free Yang-Mills, $\mathcal{N}=4$ SYM
- ❖ Understand the branch cut contributions
- ❖ Finite N ?
- ❖ ...

Thank you for the attention!