

Conserved spinning fields in conformal field theories

$$\partial_\mu \langle T_{\mu\nu} \dots \rangle = 0$$

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based on arXiv:1311.4546 and ongoing work

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Q: to what extent the full set of Ward identities constraint correlators of conserved currents?

- conformal Ward identities
invariance under the action of conformal group

$$\sum_i \mathcal{L}^i \langle S_{\mu\dots}(x_1) \dots \rangle = 0$$

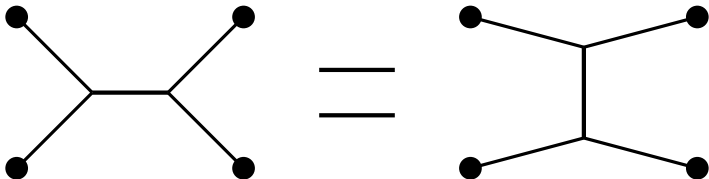
- diffeomorphism Ward identities
conservation of stress-energy tensor

$$\partial_\mu \langle S_{\mu\dots}(x_1) \dots \rangle = 0$$

Conformal symmetry or conservation can be solved independently, but so far not together

Motivation

- conformal bootstrap: 4pt function of currents and $T_{\mu\nu}$



- new formalism for CFT correlators
connection between CFT correlation functions and scattering amplitudes

Embedding formalism

- conformal group $SO(d+1, 1)$ of \mathbb{R}^d acts linearly in the embedding space $\mathbb{R}^{d+1,1}$

- space of the light-like rays $p^2 = 0$ in $\mathbb{R}^{d+1,1}$ is mapped into original \mathbb{R}^d
- auxiliary variables z^μ “take care” of spin indexes

$$\langle \mathcal{O}_{\mu\dots}(x_1) \dots \mathcal{O}_{\dots\nu}(x_n) \rangle \leftrightarrow \mathcal{P}_{\mu\dots\nu}(p_1, \dots, p_n)$$

$$P(z_1, \dots, z_n, p_1, \dots, p_n) = z_1^\mu \dots z_n^\nu \mathcal{P}_{\mu\dots\nu}(p_1, \dots, p_n)$$

$P(z, p)$ must obey certain symmetries; all such $P(z, p)$ can be built as polynomials of $H_{ij}(z, p)$ and $V_{i,jk}(z, p)$ with coefficients depending on cross-ratios

- conformal correlators $P(z, p)$ in \mathbb{R}^d are in one to one correspondence with scattering amplitudes in \mathbb{R}^{d+1}
non-conserved operators / scattering of massive particles

Solving constraints explicitly for 4pt function

- four point function of currents is parametrized by 43 functions of cross-ratios

$$\langle JJJJ \rangle = \sum_I^{43} f_I(u, v) \mathcal{P}^I(z, x)$$

$$\langle \mathcal{O}^d JJJ \rangle = \sum_{\mathcal{I}}^{14} f_{\mathcal{I}}(u, v) \mathcal{P}^{\mathcal{I}}(z, x)$$

- conservation is a system of first order linear PDE

$$\left(A_I^{\mathcal{I}}(u, v) + B_I^{\mathcal{I}}(u, v) \frac{\partial}{\partial u} + C_I^{\mathcal{I}}(u, v) \frac{\partial}{\partial v} \right) f^{\mathcal{I}}(u, v) = 0$$

matrices A, B, C are $4 \times 14 \times 43$

- $\mathbb{Z}_2 \times \mathbb{Z}_2$ permutations that respect cross-ratios

$$\Lambda_J^I(u, v) f^J(u, v) = f^I(u, v)$$

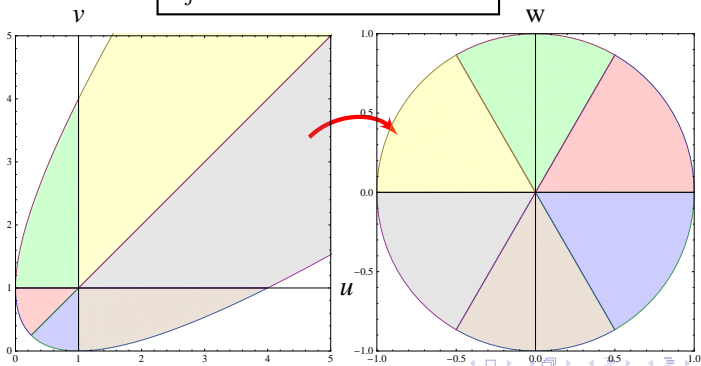
Convenient coordinates for crossing-symmetry

- bose-symmetry acts on ϕ in a simple way but leaves r invariant

$$S_{12} : \phi \rightarrow -\phi \quad S_{13} : \phi \rightarrow 2\pi/3 - \phi$$

$$u = \left| \frac{1+\lambda w}{1+w} \right|^2 \quad v = \left| \frac{1+\bar{\lambda}w}{1+w} \right|^2 \quad w = r e^{i\phi}, \quad \lambda = e^{2\pi i/3}$$

$$S_j^i(r, \phi) g^j(r, \phi) = g^i(r, \phi')$$



Cauchy problem and counting of DOF

- choosing one coordinate $t(u, v)$ as “time”
“time” derivatives of only 12 f^I out of 19 enter conservation equations

$$f^I = \underbrace{(h_1, \dots, h_{12})}_{\text{constrained}}, \underbrace{(g_1, \dots, g_7)}_{\text{unconstrained}}$$

12 equations

$$\partial_t h = \dots h + \dots \partial_x h + \dots g + \dots \partial_x g$$

2 constraints of first class

$$\dots h + \dots \partial_x h = 0$$

- crossing symmetry S_3 permutes g_i independently
 $(u, v) \rightarrow (v, u), (u/v, 1/v)$

$$S_j^i(u, v) g^j(u, v) = g^i(u', v')$$

Conformal bootstrap for four equivalent operators

split into “unconstrained” and “dependent” DOF
is not unique

$$f^I = \underbrace{(h_1, \dots, h_k)}_{\text{constrained}}, \underbrace{(g_1, \dots, g_\ell)}_{\text{unconstrained}} \Rightarrow h[g]$$

- g_i satisfy crossing-symmetry in the “bulk” (all r)

$$S_j^i(r, \phi) g^j(r, \phi) = g^i(r, \phi')$$

- h_α satisfy crossing-symmetry at the boundary $r = r_0$

$$S_\beta^\alpha(r_0, \phi) h^\beta(r_0, \phi) = h^\alpha(r_0, \phi')$$

subject to other boundary constraints

Conjecture: boundary data is irrelevant for bootstrap, i.e. if correlation function admits conformal block decomposition g fix h

Regularity and redundancy of “boundary” data

Conjecture: regular behavior at the boundary unambiguously fixes “boundary” data

ongoing work with P. Kravchuk

- example: $\langle J\mathcal{O}\mathcal{O}\mathcal{O} \rangle$

two functions, one equation

$$z = x + iy$$

$$g_-(x, y) = \frac{i}{y} \left(f(x) + \int_0^y \partial_x g_+(x, y) y dy \right)$$

regularity requires $f(x)$ to vanish

- $\langle JJJJ \rangle$: global solution can be obtained with an appropriate choice of coordinates

global solution is required to fix all boundary conditions

Bootstrap: punchline

- crossing symmetry condition applies to a small number of “independent” functions g_i ; all other DOF are ambiguously fixed
- the choice of “independent” DOF is not unique

correlator	d=3	d=4	d=5	d \geq 6
$\langle JJJJ \rangle$	5	7	7	7
$\langle TTTT \rangle$	5	22	28	29

number of “independent” functions/families of conformal blocks relevant for corresponding parity-even 4pt functions in various dimensions

Finding general solution

- zero curvature equation $\partial_\mu A_\nu - \partial_\nu A_\mu = 0 \Rightarrow A_\mu = \partial_\mu \alpha$

$$\text{example } \langle J\mathcal{O}\mathcal{O}\mathcal{O} \rangle: \quad g_\pm = \frac{\pm \partial_\pm \alpha}{z - \bar{z}}$$

- special solution/conformal blocks to satisfy Ward identities for charged operators

$$\partial_\mu \langle J_\mu(x) \dots \rangle = \delta(x - y) \dots$$

- system of differential equations $Df = 0$ can be solved in terms of compatibility operator \mathcal{D} , $f = \mathcal{D}\alpha$,

$$D\mathcal{D} = 0$$

there is an algorithmic way to construct \mathcal{D} ; for $\langle JJJJ \rangle$ in $d = 3$ there are 5 functions α , as expected;

New formalism?

Relation between CFT correlators and scattering amplitudes

- scattering amplitudes for massless particles can be constructed in terms of V, H
- choice of compatibility operator \mathcal{D} is not unique
 - ▶ for $\langle J_\mu \dots \rangle$ solution can be expressed as $\partial_\nu \langle F_{\mu\nu} \dots \rangle$
 - ▶ more generally $\tilde{\mathcal{D}}$ can be built of $z_{[m} p_n \frac{\partial}{\partial p^k]}$

$$D\tilde{\mathcal{D}} = 0$$

no clear understanding yet how to construct *all* conformal correlators that include conserved fields

Conclusions

- conservation significantly reduces number of “independent” DOF which drastically simplifies (numerical) bootstrap; “boundary” data is fixed by regularity
- all set of Ward identities (including conservation) can be solved explicitly in terms of differential operators acting on independent functions; such a representation is not unique
- at this moment there is no universal mathematically elegant formalism to express conformal correlators in terms of unrestricted DOF that would make a connection with scattering amplitudes