

with Wenliang Li and Massimo Taronna [1406.2335]

On partially massless gravity

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Graviton: **massless** spin-② particle

$$S = \int d^D x - \frac{1}{2} \partial_\lambda h_{\mu\nu} \partial^\lambda h^{\mu\nu} + \partial_\mu h_{\nu\lambda} \partial^\nu h^{\mu\lambda} - \partial_\mu h^{\mu\nu} \partial_\nu h + \frac{1}{2} \partial_\lambda h \partial^\lambda h$$

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massive spin-**2** particle

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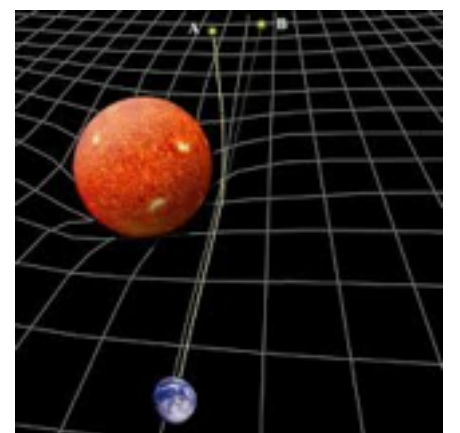
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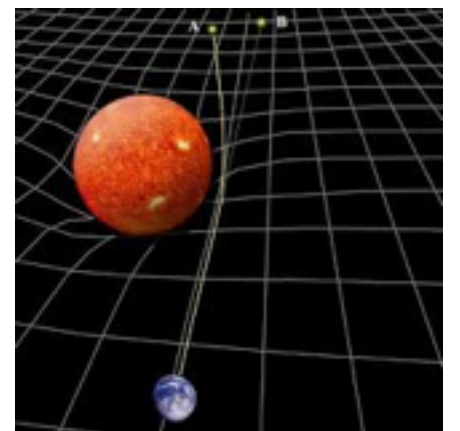
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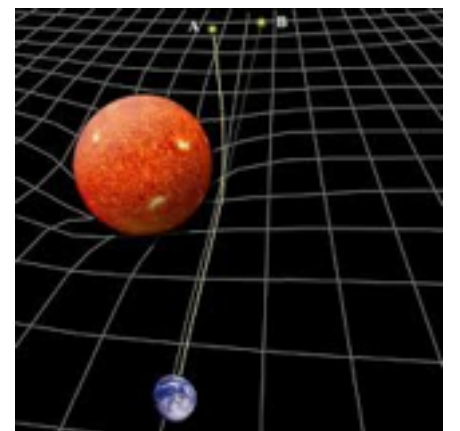
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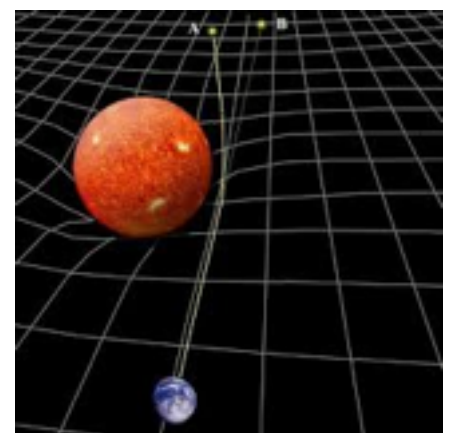
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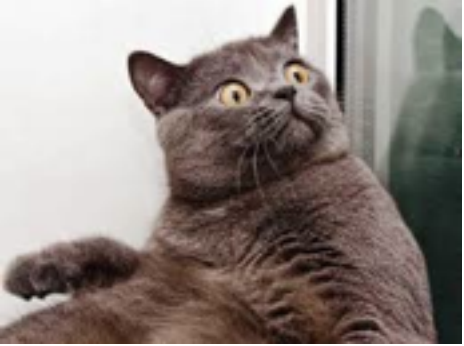
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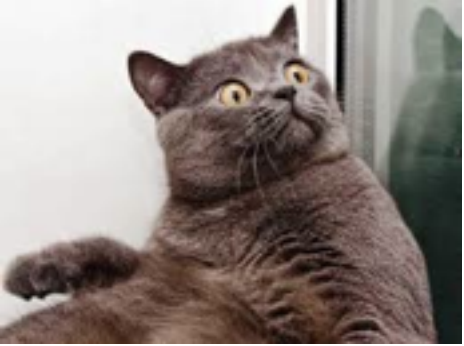
- Vainshtain mechanism

non-linearity removes the discontinuity



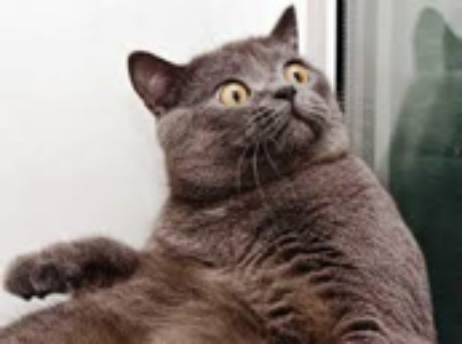


Surprises in constant curvature background



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- **AdS** : **NO** DVZ discontinuity **②** $\rightarrow$ ② + **①**



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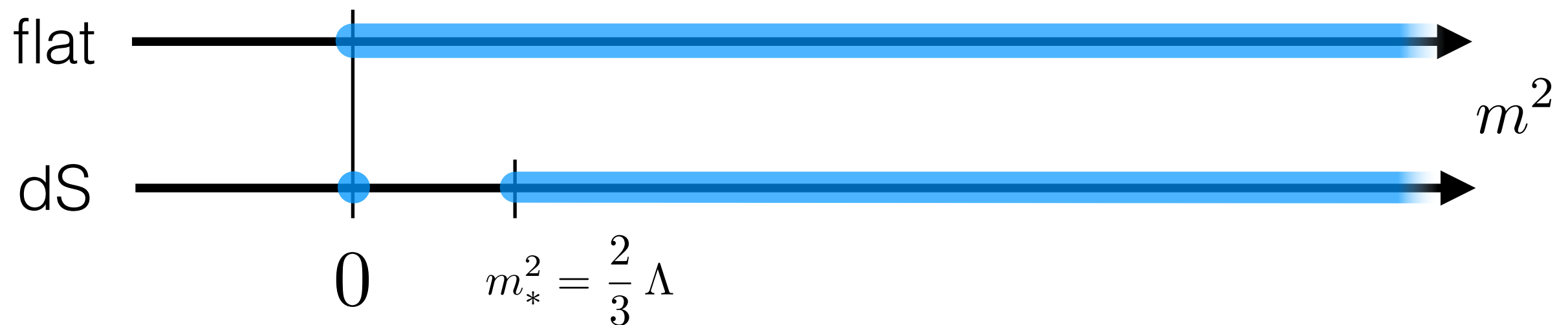




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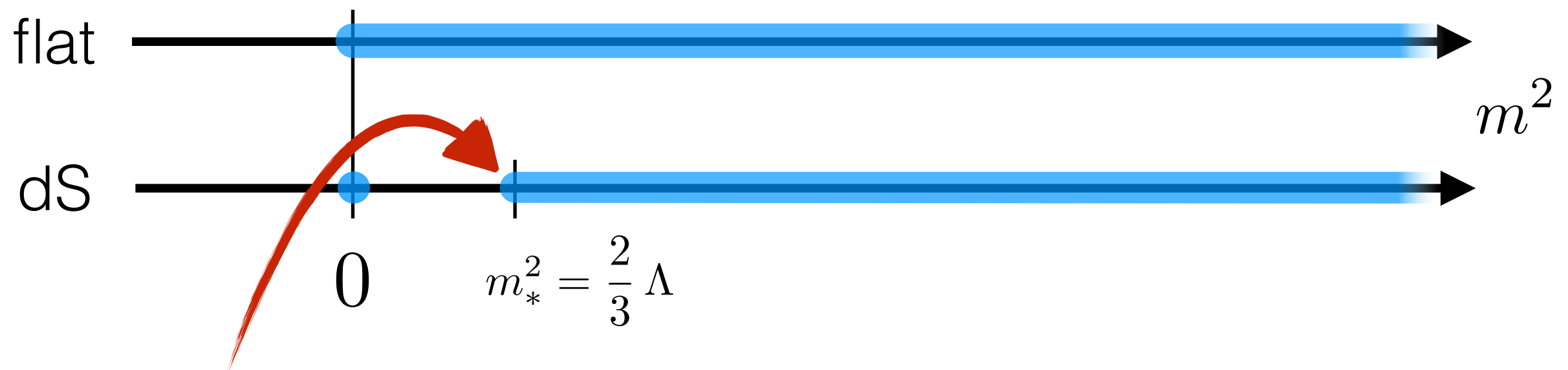




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the end point is special!

$$\textcircled{2} \rightarrow \textcircled{2} + \textcircled{0}$$

Partially massless spin-2 field

Deser, Waldron

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- Smaller number of DoF “ $\textcircled{2} = \textcircled{2} + \textcircled{1}$ ”

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- Does this spectrum appear in some theories?

Massive gravity / Bi-gravity

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What happens when $m^2 \rightarrow m_*^2$?

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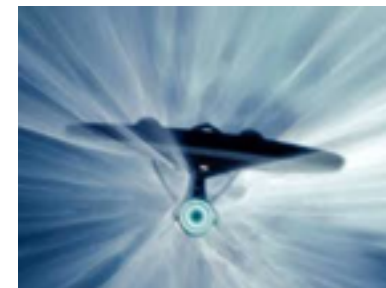
② $\rightarrow$ ② + ① decouples from the theory?

① is responsible for

- DVZ discontinuity
(it is cured by Vainstein mechanism)

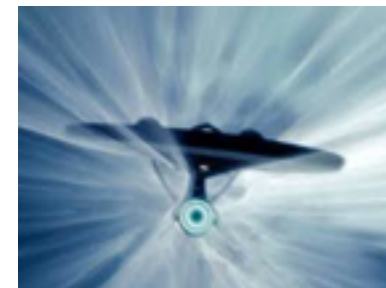
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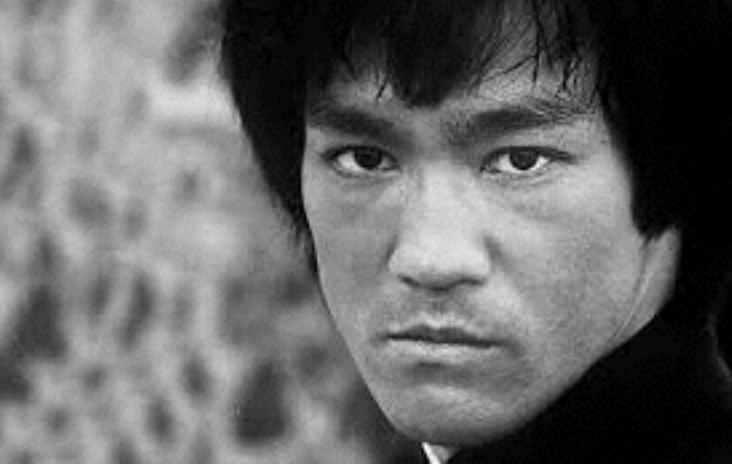
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Does it decouple from Massive gravity / Bi-gravity
when $m^2 \rightarrow m_*^2$?

maybe yes vs maybe no



a definite answer to this question



a definite answer to this question
from
cubic interaction analysis





a definite answer to this question
from
cubic interaction analysis & a little more



Let us reformulate the task:

Find out **all possible interacting theories**
for partially massless ② and massless ② !

General properties of an action $\mathcal{S}[\chi_i]$
invariant under **gauge** transf. $\delta_\varepsilon \chi_i$

Berends, Burger, van Dam;
Barnich, Henneaux

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strategy 1

expansion in χ_i :

$$S = S^{[2]} + S^{[3]} + \dots,$$

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Admissibility condition Vasiliev



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$\chi_1 = h$ massless ②

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Zinoviev

- Two-derivative interactions



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- $\mathcal{L}_3$ exists **only in 4d**



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- α, β, γ : coupling constants



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$$\delta^{[1]} h_{\mu\nu} = \alpha (2 \bar{\nabla}_{(\mu} \bar{\xi}^{\rho} h_{\nu)\rho} + \bar{\xi}^{\rho} \bar{\nabla}_{\rho} h_{\mu\nu}) - 2 \beta \partial^{\rho} \bar{\alpha} (2 \bar{\nabla}_{(\mu} \varphi_{\nu)\rho} - \bar{\nabla}_{\rho} \varphi_{\mu\nu})$$



$$\delta^{[0]} S^{[2]} = 0$$

$$\delta^{[0]} S^{[3]} + \delta^{[1]} S^{[2]} = 0$$

$$S^{[2]} = \int h \partial^2 h + \sigma \varphi (\partial^2 + \frac{2}{3} \Lambda) \varphi \quad S^{[3]} = \int \alpha \mathcal{L}_1(h, h, h) + \beta \mathcal{L}_2(h, \varphi, \varphi) + \gamma \mathcal{L}_3(\varphi, \varphi, \varphi)$$

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❖ so far, coupling constants α, β, γ are not determined



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❖ Global sym: $\begin{cases} SO(2, 4) & \text{for } \sigma \Lambda < 0 \\ SO(1, 5) & \text{for } \sigma \Lambda > 0 \end{cases}$



Admissibility condition

$$\delta_{\bar{\varepsilon}}^{[1]} \delta_{\bar{\eta}}^{[1]} - \delta_{\bar{\eta}}^{[1]} \delta_{\bar{\varepsilon}}^{[1]} = \delta_{[\bar{\eta}, \bar{\varepsilon}]}^{[1]} + \delta_{[\bar{\eta}, \bar{\varepsilon}]^{[1]}}^{[0]} + C_{ij}^{[0]}(\bar{\eta}, \bar{\varepsilon}) \frac{\delta S^{[2]}}{\delta \chi_i} \frac{\delta}{\delta \chi_j}$$



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- When $\sigma = -1$, we recover **Conformal Gravity**